

Kittel

$$4.1. (a) T = \frac{1}{2} M v^2 = \frac{1}{2} M \left(\frac{du_s}{dt} \right)^2 \text{ per atom.}$$

In the context of only considering nearest neighbor interactions, we use the harmonic oscillator potential.

$$U = \frac{1}{2} k x^2 = \frac{1}{2} C (u_s - u_{s+1})^2 \text{ between } s \text{ and } s+1.$$

$$\Rightarrow E = \frac{1}{2} M \sum_s \left(\frac{du_s}{dt} \right)^2 + \frac{1}{2} C \sum_s (u_s - u_{s+1})^2$$

In the second sum, we don't have to sum both neighbors of an atom (i.e., $(u_s - u_{s+1})^2 + (u_s - u_{s-1})^2$) because the potential is stored in the "spring" not in the atom. In a linear lattice each atom has on average one spring connected to its neighbor on the right.

$$(b). \left\langle \frac{1}{2} M \left(\frac{du_s}{dt} \right)^2 \right\rangle = \frac{1}{2} M u^2 \langle \cos^2(\omega t - s k a) \rangle \\ = \frac{1}{4} M u^2$$

$$\left\langle \frac{1}{2} C (u_s - u_{s+1})^2 \right\rangle = \frac{1}{2} C \left\langle u^2 \left[\cos(\omega t - s k a) - \cos(\omega t - (s+1) k a) \right]^2 \right\rangle \\ = \frac{1}{2} C u^2 \left\langle \left[\cos^2(X) + \cos^2(X - k a) - 2 \cos(X) \cos(X - k a) \right] \right\rangle \\ = \frac{1}{2} C u^2 \left[1 - 2 \langle \cos X \cos(X - k a) \rangle \right], \quad X = \omega t - s k a \\ = \frac{1}{2} C u^2 \left[1 - 2 \left[\frac{1}{2} \langle \cos(2X - k a) \rangle + \frac{1}{2} \langle \cos k a \rangle \right] \right] \\ = \frac{1}{2} C u^2 [1 - \cos k a]$$